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Damping Capacity of Aluminum 6061-Indium Alloys

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**Damping Capacity of Aluminum 6061-  
Indium Alloys**

by

C.R. Wong

O. Diehm\*

D.C. Van Aken\*

\*Department Materials Science and Engineering

University of Michigan

Ann Arbor, MI 48109



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# CONTENTS

|                                 |    |
|---------------------------------|----|
| ABSTRACT.....                   | 1  |
| ADMINISTRATIVE INFORMATION..... | 1  |
| INTRODUCTION.....               | 1  |
| NIELSEN MODEL.....              | 2  |
| EXPERIMENTAL PROCEDURE.....     | 7  |
| RESULTS.....                    | 8  |
| DISCUSSION.....                 | 10 |
| CONCLUSIONS.....                | 11 |
| ACKNOWLEDGMENTS.....            | 12 |
| REFERENCES.....                 | 13 |

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| A-1                  |                                     |



## FIGURES

|           |   |    |
|-----------|---|----|
| Figure 1: | Single Cantilever Clamping Arrangement of the<br>Dynamic Mechanical Thermal Analyzer.....   | 14 |
| Figure 2: | Microstructures of 6061-T6 Aluminum with<br>1.4 - 5.2 Volume Percent Indium.....  | 15 |
| Figure 3: | Precipitates in 6061-T6 Aluminum with<br>0 - 5.2 Volume Percent Indium.....   | 16 |
| Figure 4: | Values of the Loss Factor and Storage Modulus of the<br>Monolithic Materials Measured from 20 to 100°C<br>at 0.1 Hz.....                            | 17 |
| Figure 5: | Values of the Loss Factor and Storage Modulus of<br>6061 Aluminum with 0 to 2.7 Volume Percent Indium<br>Measured from 20 to 100°C at 0.1 Hz.....   | 18 |
| Figure 6: | Values of the Loss Factor and Storage Modulus of<br>6061 Aluminum with 0 to 5.2 Volume Percent Indium<br>Measured from 20 to 100°C at 0.1 Hz.....   | 19 |
| Figure 7: | Values of the Loss Factor and Storage Modulus of<br>6061 Aluminum with 0 to 2.7 Volume Percent Indium<br>Calculated Using Nielsens Model.....       | 20 |
| Figure 8: | Comparison of Measured and Calculated Values of<br>the Loss factor and Storage Modulus of 6061 Aluminum<br>with 0 to 2.7 Volume Percent Indium..... | 21 |
| Figure 9: | Comparison of Measured and Calculated Values of<br>the Loss Factor and Storage Modulus of 6061 Aluminum<br>with 5.2 Volume Percent Indium.....      | 22 |

## TABLE

|          |   |    |
|----------|---|----|
| Table 1: | Chemical Composition of Al-6061-In-T6 Alloys..... | 13 |
|----------|---|----|

## ABSTRACT

Type 6061 aluminum alloys containing between 0 and 5.2 volume percent indium and pure indium samples were fabricated. Each sample was characterized by metallographic and analytical electron microscopy and the damping capacity and storage modulus were measured. The model proposed by L. G. Nielsen was used to calculate the damping capacity and storage modulus of the alloys using the damping capacity and storage modulus of pure indium and 6061 aluminum. The damping capacity of the Al-6061-In-T6 alloys were higher than the Al-6061-T6 alloy and increased with increasing indium content. The Nielsen model gave a good first approximation of the damping capacity and storage modulus of the alloys.

## ADMINISTRATIVE INFORMATION

This report was prepared under the Quiet Alloys program, part of the Functional Materials Block Program, under the sponsorship of Mr. Ivan Caplan, David Taylor Research Center (DTRC Code 0115). Work was performed at the David Taylor Research Center and the Department of Materials Science and Engineering, University of Michigan, Ann Arbor. The work was supervised by Dr. O. P. Arora, DTRC Code 2812, under Program Element 62234N, Task Area RS34S94, Work Unit 1-2812-949. This report satisfies FY89 Milestone 94SR1/6.

## INTRODUCTION

An important characteristic of a structural material is its damping capacity. While metallic materials exhibit adequate stiffness for structural use, the damping capacity may be quite low, having a typical loss factor on the order of  $10^{-4}$ . In contrast, polymeric materials exhibit very high damping, with loss factors on the order of one, but rather low stiffness. Their stiffness can be increased with the use of fillers and fibers but the resultant resin matrix composites exhibit lower damping properties, with loss factors on the order of  $10^{-2}$ . Attempts made to improve the damping response of the resin matrix composite by adding rubber did not result in significant improvements [1]. It was shown that synergistic effects from interactions

between the rubber and the resin were responsible for the lower than expected damping behavior.

In the case of metal matrix composites, work by Ray, Kinra, Rawal and Misra has shown that the damping of aluminum alloy 6061 is increased by the addition of graphite fibers [2]. However, the increase in damping was low considering the high volume fraction (0.34) of graphite. Recent work by Diehm, Wong and Van Aken has shown that the addition of a viscoelastic inclusion (indium) to pure aluminum will produce high damping materials [3], but it was uncertain whether the principal damping resulted from the matrix or the inclusion since both have high damping capacities.

In this investigation the addition of indium, an elastically soft second phase particle, to 6061 T6 aluminum, a stiff matrix, was examined in order to differentiate between inclusion and matrix damping. Additionally the model to predict the stiffness of composite materials proposed by L. G. Nielsen [4,5] was evaluated for its ability to predict the damping capacity of composite materials. The dynamic properties of pure (99.99%) indium and 6061 T6 aluminum were determined. The dynamic properties of the composite were calculated using the values of the monolithic materials in the Nielsen model and directly compared with the experimental results.

#### NIELSEN MODEL

The model developed by Nielsen [4] predicts the complex modulus of isotropic two phase materials with arbitrary phase geometry. It is based on a continuum mechanics composite sphere assemblage model but is semi-empirical. The model assumes that the alloy is isotropic, strained only in the elastic range, and is phase symmetric, that is both the matrix and second phase geometries are identical at equal respective volume concentrations. Equations

1-4 below, from Nielsen's model [5], can be used to calculate Young's modulus of the alloy,  $E_y$ , using the Young's moduli of the matrix,  $E_y^s$ , and second

phase,  $E_y^i$ , and the volume concentration,  $c$ . The volume concentration =  $\frac{V^i}{(V^s + V^i)}$

where  $V^i$  and  $V^s$  are the volumes of the second phase and matrix respectively.

$$E_y = eE_y^s \quad \text{eq.1}$$

where  $e$  is the relative Young's modulus of the alloy.

$$e = \frac{n + \gamma + \gamma c(n - 1)}{n + \gamma - c(n - 1)} \quad \text{eq.2}$$

where  $n$  is the relative stiffness and  $\gamma$  is the shape function.

$$n = \frac{E_y^i}{E_y^s} \quad \text{eq.3}$$

$$\gamma = \frac{1}{2} \left\{ \rho[1 - c(1 - n)] + \sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)} \right\} \quad \text{eq.4}$$

where  $\rho$  is the shape factor which is dependent on the morphology of the composite.

The complex modulus of the matrix,  $E^i$ , and second phase,  $E^s$ , is defined as follows.

$$E^s = a^s + b^s i \text{ and } E^i = a^i + b^i i \quad \text{eq.5}$$

where  $a$  and  $b$  are the storage and loss modulus respectively and the superscripts  $s$  and  $i$  refer to the matrix and second phase respectively. The conversion from Young's modulus equations to complex modulus equations is accomplished with the use of the correspondence principle. The complex moduli from equation 5 are substituted for the Young's moduli in equations 1 and 3 and the real and imaginary parts are separated. Starting with equation 3:

$$n = \frac{E^i}{E^s} = \frac{a^i + ib^i}{a^s + ib^s} = \frac{(a^i + ib^i)(a^s - ib^s)}{(a^s + ib^s)(a^s - ib^s)} = \frac{a^i a^s + b^i b^s}{(a^s)^2 + (b^s)^2} + \frac{(a^s b^i - a^i b^s)i}{(a^s)^2 + (b^s)^2}$$

$$\text{Let } n = A + Bi \text{ where } A = \frac{a^i a^s + b^i b^s}{(a^s)^2 + (b^s)^2} \text{ and } B = \frac{a^s b^i - a^i b^s}{(a^s)^2 + (b^s)^2} \quad \text{eq. 6}$$

Now recalling equation 4

$$\gamma = \frac{1}{2} \left\{ \rho[1 - c(1 - n)] + \sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)} \right\}$$

Upon substitution of equation 6 the first part of equation 4 becomes

$$\rho[1 - c(1 - n)] = \rho - \rho c + \rho c n = \rho - \rho c + \rho c A + \rho c B i \quad \text{eq. 7}$$

The second part of equation 4 is  $\sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)}$

$$\begin{aligned} [1 - c(1 - n)]^2 &= 1 - 2c(1 - n) + c^2(1 - n)^2 \\ &= 1 - 2c + 2cn + c^2 - 2c^2n + c^2n^2 \\ &= 1 - 2c + c^2 + (2c - 2c^2)n + c^2n^2 \end{aligned}$$

$$\text{since } n^2 = (A + iB)(A + iB) = A^2 - B^2 + 2ABi$$

$$\begin{aligned} \text{then } [1 - c(1 - n)]^2 &= (1 - 2c + c^2) + (2c - 2c^2)A + c^2(A^2 - B^2) \\ &\quad + i[(2c - 2c^2)B + 2c^2AB] \end{aligned}$$

therefore  $\sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)}$

$$\begin{aligned}
&= \left\{ \rho^2[1 - 2c + c^2 + (2c - 2c^2)A + c^2(A^2 - B^2)] \right. \\
&\quad \left. + i\rho^2[(2c - 2c^2)B + 2c^2AB]i + 4A(1 - \rho) + 4B(1 - \rho) \right\}^{\frac{1}{2}} \\
&= \left\{ \rho^2[1 - 2c + c^2 + 2c(1 - c)A + c^2(A^2 - B^2)] \right. \\
&\quad \left. + 4A(1 - \rho) + i[\rho^2 2c(1 - c)B + 2c^2AB\rho^2 + 4B(1 - \rho)] \right\}^{\frac{1}{2}}
\end{aligned}$$

$$\text{Let } \sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)} = [\alpha + \beta i]^{\frac{1}{2}} \quad \text{eq.8}$$

$$\text{where } \alpha = \rho^2[(c - 1)^2 - 2c(c - 1)A + c^2(A^2 - B^2)] + 4A(1 - \rho) \quad \text{eq.9}$$

$$\text{and } \beta = \rho^2 2c(1 - c)B + 2c^2AB\rho^2 + 4B(1 - \rho) \quad \text{eq.10}$$

In order to find the square root the coordinates are changed.

$$r = (\alpha^2 + \beta^2)^{\frac{1}{2}} \quad \text{eq.11}$$

$$\theta = \arctan \left( \sqrt{\frac{\beta}{\alpha}} \right) \quad \text{eq.12}$$

substituting equations 11 and 12 into equation 8 results in

$$\sqrt{\rho^2[1 - c(1 - n)]^2 + 4n(1 - \rho)} = r^{1/2}[\cos(\theta/2) + i\sin(\theta/2)] = r^{1/2}e^{i\theta/2} \quad \text{eq.13}$$

Combining equations 7 and 13 gives the complex shape function,  $\gamma^*$ .

$$\gamma^* = \frac{1}{2} \left\{ \rho[1 - c(1 - A)] + \rho c B i + r^{1/2} e^{i\theta/2} \right\} \quad \text{eq.14}$$

Substituting the complex values of  $\gamma^*$  from equation 14 and the complex values of  $n$  from equation 6 into equation 2 gives the complex relative modulus,  $e^*$ .

$$\begin{aligned} e^* &= \frac{n + \gamma^* + \gamma^* c(n - 1)}{n + \gamma^* - c(n - 1)} = \frac{n + \gamma^* + c n \gamma^* - \gamma^* c}{n + \gamma^* - c n + c} \\ &= \frac{A + \text{Re}(\gamma^*) - c \text{Re}(\gamma^*) + c[A \text{Re}(\gamma^*) - B \text{Im}(\gamma^*)]}{(A + \text{Re}(\gamma^*) - cA + c) + (B + \text{Im}(\gamma^*) - cB)} \\ &\quad + \frac{i \{ B + \text{Im}(\gamma^*) - c \text{Im}(\gamma^*) + c[A \text{Im}(\gamma^*) + B \text{Re}(\gamma^*)] \}}{(A + \text{Re}(\gamma^*) - cA + c) + (B + \text{Im}(\gamma^*) - cB)} \end{aligned} \quad \text{eq.15}$$

$$\text{Let } \xi = A + \text{Re}(\gamma^*) - c \text{Re}(\gamma^*) + c[A \text{Re}(\gamma^*) - B \text{Im}(\gamma^*)] \quad \text{eq.16}$$

$$\text{and } \eta = B + \text{Im}(\gamma^*) - c \text{Im}(\gamma^*) + c[A \text{Im}(\gamma^*) + B \text{Re}(\gamma^*)] \quad \text{eq.17}$$

and substitute into equation 15.

$$\begin{aligned} e^* &= \frac{\xi + i\eta}{(A + \text{Re}(\gamma^*) - cA + c) + (B + \text{Im}(\gamma^*) - cB)i} \\ &= \frac{(\xi + i\eta)[(A + \text{Re}(\gamma^*) - cA + c) - (B + \text{Im}(\gamma^*) - cB)]}{(A + \text{Re}(\gamma^*) - cA + c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \\ &= \frac{\xi(A + \text{Re}(\gamma^*) - cA + c) + \eta(B + \text{Im}(\gamma^*) - cB)}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \\ &\quad + \frac{i\eta(A + \text{Re}(\gamma^*) - cA + c) - \xi(B + \text{Im}(\gamma^*) - cB)}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \end{aligned} \quad \text{eq.18}$$

Finally the complex modulus of the alloy is found by combining

equations 1, 5 and 18.

$$E^i = e^* E^s = \text{Re}(e^*) a^s - \text{Im}(e^*) + i[\text{Im}(e^*) + \text{Re}(e^*) b^s]$$

$$\begin{aligned} &= \frac{a^s \left\{ \xi(A + \text{Re}(\gamma^*) - cA + c) + \eta(B + \text{Im}(\gamma^*) - cB) \right\}}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \\ &- \frac{b^s \left\{ \eta(A + \text{Re}(\gamma^*) - cA + c) - \xi(B + \text{Im}(\gamma^*) - cB) \right\}}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \\ &+ i \left\{ \frac{a^s \left\{ \eta(A + \text{Re}(\gamma^*) - cA + c) - \xi(B + \text{Im}(\gamma^*) - cB) \right\}}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \right. \\ &\left. + \frac{b^s \left\{ \xi(A + \text{Re}(\gamma^*) - cA + c) + \eta(B + \text{Im}(\gamma^*) - cB) \right\}}{(A + \text{Re}(\gamma^*) - cA - c)^2 + (B + \text{Im}(\gamma^*) - cB)^2} \right\} \end{aligned} \quad \text{eq.19}$$

Where the real part of equation 19 is the storage modulus of the composite and the imaginary part of equation 19 is the loss modulus.

#### EXPERIMENTAL PROCEDURE

Aluminum 6061 alloys with additions of 0 to 13 weight percent indium were prepared by plasma arc-melting. The starting alloys were pure indium (99.99%) and 6061 alloy. The chemical composition of the alloys were determined by wet-chemistry. The volume fraction of indium was calculated using the weight fraction and density of each alloy by assuming complete immiscibility between aluminum and indium. The arc-melted ingot was then reduced 60 to 80% in thickness, by repeatedly cold-rolling 20 to 30% and annealing, to produce a flat sample with a nominal thickness of 1.5 mm. The alloys were given a T6 temper consisting of solution treatment at 532 °C (990 °F) and aging at 193 °C (380 °F) for 7 hours. Samples of pure indium were likewise plasma arc-melted and rolled.

Each sample was characterized by metallographic and analytical electron microscopy. Electron microscopy studies were performed at the University of Michigan Electron Microbeam Analysis Laboratory. Thin foils for transmission electron microscopy were prepared by twin jet electropolishing in a solution of 20% nitric acid (by volume) and methanol.

The damping capacity and modulus of the samples were measured with a Polymer Laboratories Dynamic Mechanical Thermal Analyzer (DMTA). The DMTA uses a fixed-guided cantilevered arrangement where the left clamp holds the sample to a stationary frame while the right clamp attaches the sample to the drive shaft as illustrated in Fig. 1. A small sinusoidal mechanical stress is applied to the cantilevered sample and the resulting sinusoidal strain is measured with a noncontacting eddy current transducer. Comparison of the amplitude of the stress,  $\sigma$ , and strain,  $\epsilon$ , signals yields the storage modulus,  $a$ , and the phase lag of strain behind the stress gives the phase angle,  $\delta$ . The complex modulus,  $E$ , and loss modulus,  $b$ , are calculated using the following equation:

$$\frac{\sigma}{\epsilon} = a(1 + i \tan \delta) = E = a + ib \quad \text{eq.20}$$

where  $\tan \delta$  is the loss factor. The frequency of the vibrations was cycled between 0.1, 1 and 10 Hz while the temperature was increased one degree C per minute from 20 °C (68 °F) to 100 °C (212 °F). Each sample was measured at least twice to check measurement consistency.

## RESULTS

The measured chemical composition and the calculated volume fraction of indium are presented in table 1. The volume percent varied from 0 to 5.2. The microstructures of the indium containing alloys are shown in Fig. 2. A

uniform dispersion of indium particles was found in all the samples with the individual areas of indium increasing in size and number with the increase in volume percent. The micrographs show the indium phase to be roughly spherical. Examination of the age-hardened matrix using transmission electron microscopy revealed that the age-hardening process was affected by the addition of indium. A typical 6061 T6 microstructure consists of a uniform distribution of Guinier-Preston Zones (GPZ) and  $\beta'$  (rod shaped  $Mg_2Si$ ) precipitates in the aluminum matrix as shown in Fig. 3a. The diffraction conditions are optimized in Figs. 3a and 3b to show the  $\beta'$  precipitates. The aged microstructures of the alloys containing 1.4, 1.7 and 5.2 volume percent indium are shown in Figs. 3b to 3d. It is apparent that the aging kinetics have been affected by the additions of indium. The general trend is that the precipitation of  $\beta'$  is inhibited and the volume fraction of second phase is reduced. Only the GPZ's are observed in the 1.7 and 5.2% alloys.

The results of the DMTA testing are shown as plots of loss factor,  $\tan\delta$ , versus the storage modulus on logarithmic axis in order to eliminate temperature and frequency measurement error from the data. As the temperature was increased from 20 °C to 100 °C the loss factor increased and the storage modulus decreased. The measurements of pure indium and the 6061 T6 alloy are shown in Fig. 4. For the temperature range tested, the storage modulus of the 6061 T6 alloy did not vary significantly from 71 GPa while the storage modulus of the indium varied from 2 GPa at room temperature to 0.9 GPa at 100 °C. It was generally observed that the storage modulus decreased and the loss factor increased with increasing addition of indium as shown in Fig. 5. The storage modulus of the sample containing 5.2 volume percent indium exhibited a more dramatic change than alloys containing less than 3.2 volume percent indium, as illustrated in Fig. 6. The loss factor of the 5.2

volume percent indium alloy at room temperature was measured to be 0.01. This was likely due to increased continuity of the indium phase. The storage modulus and loss factor were calculated with the Nielsen Model using the data from the monolithic material in equations 19 and 20 and a shape factor of one. A shape factor of one describes perfectly spherical second phase areas completely surrounded by the matrix. The results of these calculations are presented in Fig. 7. The calculated and measured values of the 0 volume percent indium alloy are constrained to be equal. Comparing the calculated values to the measured values as in Figs. 8 and 9 it is obvious that although the calculated values show the same trends as the measured values, they consistently overestimate both the measured storage modulus and the loss factor of the alloys. For the alloys containing less than 3.2 volume percent the storage modulus is only overestimated by 2% and the loss factor is overestimated by 30%. However, in the case of the 5.2 volume percent indium alloy the storage modulus was overestimated by more than 100% while the loss factor was overestimated by 60%. These results may indicate a synergistic effect such as the partitioning of alloying elements present in the 6061 material to the indium.

#### DISCUSSION

High damping aluminum alloys may be obtained by the addition of a viscoelastic inclusion. In the present case a volume fraction of at least 0.05 is required to produce an alloy with a loss factor greater than 0.01. However, there is a significant loss of stiffness associated with the addition of the indium and there appears to be a synergistic effect between the matrix and the inclusion. The aged 6061-T6 microstructure shows a decreasing precipitate density with increasing indium content and the measured loss

factors are much less than the calculated values based on the damping capacities of the monolithic samples. It is tempting to speculate that these observations are related. Indeed, the solubility of magnesium in indium is greater than 30 atomic percent at the T6 aging temperature used in this experiment [6]. Thus the low volume fraction of precipitates may be related to the partitioning of magnesium to the indium inclusions. Furthermore, The indium-magnesium inclusions may have a lower damping capacity than the pure indium. If indeed the damping of the indium inclusion is a strain dependent mechanism, such as dislocation motion, the addition of solute atoms will result in a lower loss factor for a comparable cyclic strain.

The Nielsen model failed to predict the dynamic properties of indium containing 6061 T6 alloys from the pure indium and 6061 T6 alloy properties, but did provide a good first approximation. Future modeling of this system will use the dynamic properties of monolithic indium-magnesium alloys to compensate for the synergistic effects encountered and the shape factor will be varied in an attempt to compensate for inclusions which are not perfectly spherical.

#### CONCLUSIONS

Additions of indium, an elastically soft second phase particle, to 6061 T6 aluminum, a stiff matrix, have resulted in an increased damping capacity while still maintaining the stiffness of the matrix. The measured and calculated values agree that the damping capacity increases and the storage modulus decreases with increasing indium content. The Nielsen model is a good first approximation for both the prediction of the maximum damping capacity and stiffness of a particular alloy system and the tailoring of

alloys to obtain the damping capacity and stiffness required by a given application.

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Table 1: Chemical Composition of Al-6061-In-T6 Alloys

| Calculated<br>Volume Percent<br>Indium | Measured<br>Weight Percent |           |          |         |        |      |          |
|--|----------------------------|-----------|----------|---------|--------|------|----------|
|  | Indium                     | Magnesium | Chromium | Silicon | Copper | Iron | Aluminum |
| 0.00                                   | 0.00                       | 0.77      | 0.048    | 0.71    | 0.26   | 0.23 | 98.97    |
| 0.78                                   | 2.08                       | 0.74      | 0.047    | 0.83    | 0.27   | 0.25 | 95.78    |
| 1.43                                   | 3.77                       | 0.70      | 0.046    | 0.76    | 0.26   | 0.24 | 94.22    |
| 1.67                                   | 4.37                       | 0.67      | 0.045    | 0.73    | 0.25   | 0.22 | 93.72    |
| 2.16                                   | 5.63                       | 0.70      | 0.044    | 0.75    | 0.26   | 0.22 | 92.40    |
| 2.66                                   | 6.87                       | 0.73      | 0.045    | 0.71    | 0.25   | 0.21 | 91.19    |
| 3.20                                   | 8.20                       | 0.73      | 0.041    | 0.70    | 0.28   | 0.22 | 89.83    |
| 5.16                                   | 12.80                      | 0.70      | 0.042    | 0.64    | 0.23   | 0.20 | 85.39    |

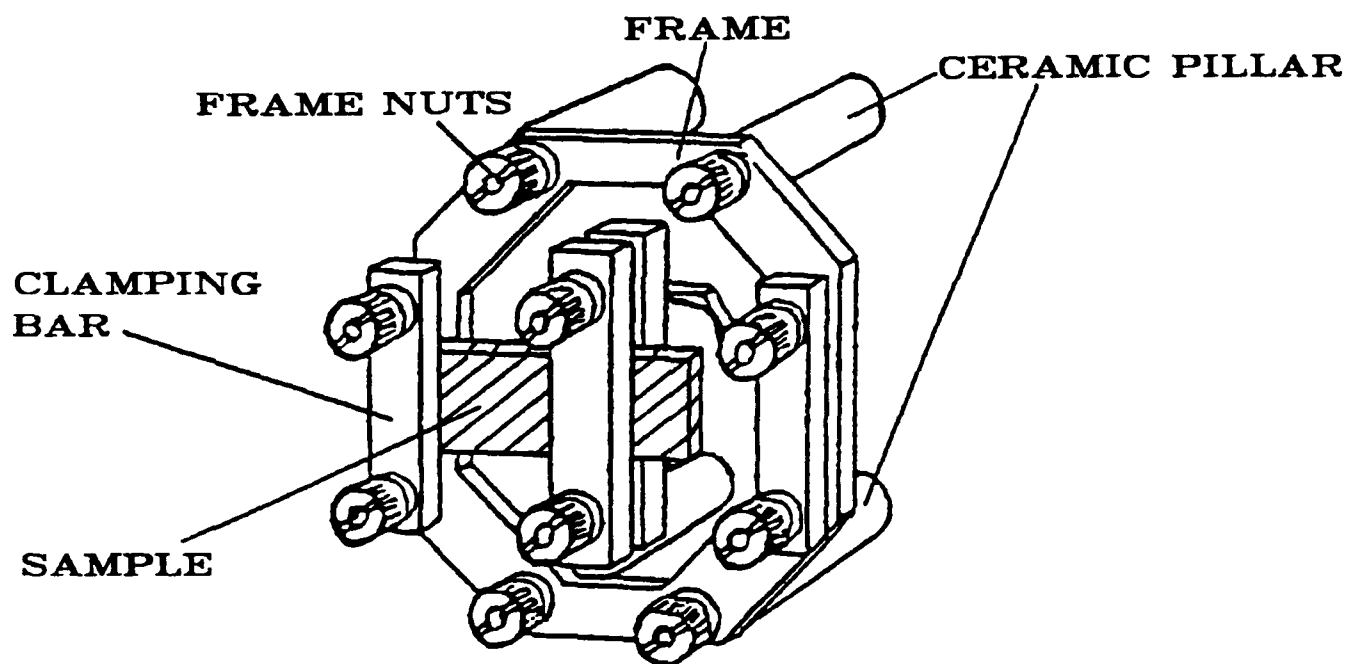


Fig. 1. Single cantilever clamping arrangement of the dynamic mechanical thermal analyzer.

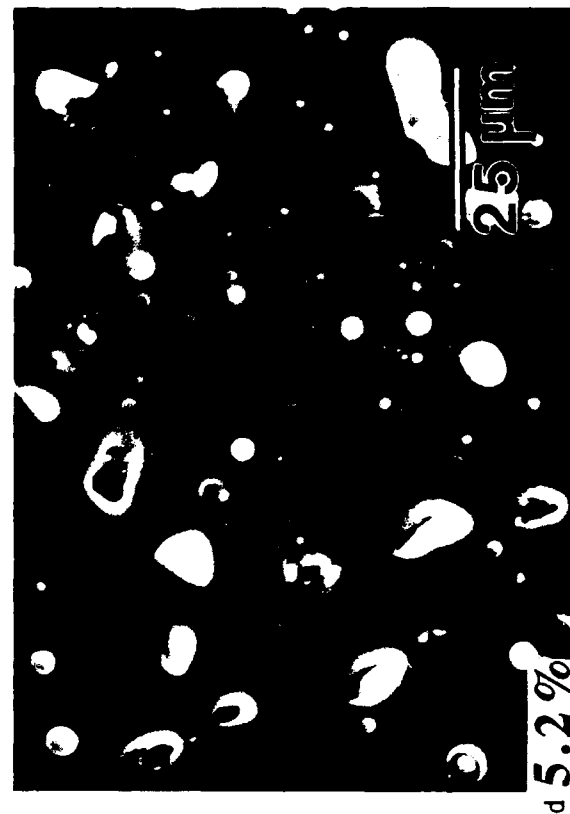


Fig. 2. Microstructures of 6061-T6 aluminum with 1.4 – 5.2 volume percent indium.

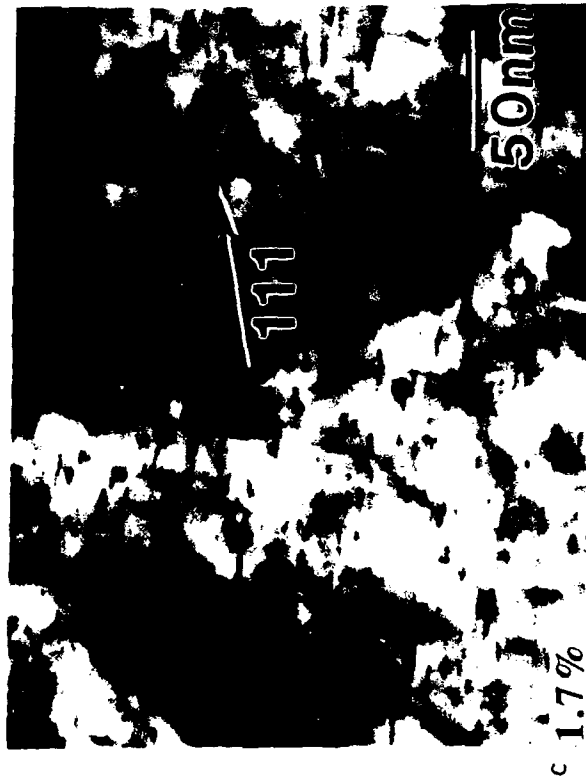


Fig 3. Precipitates in 6061-T6 aluminum with 0 – 5.2 volume percent indium.

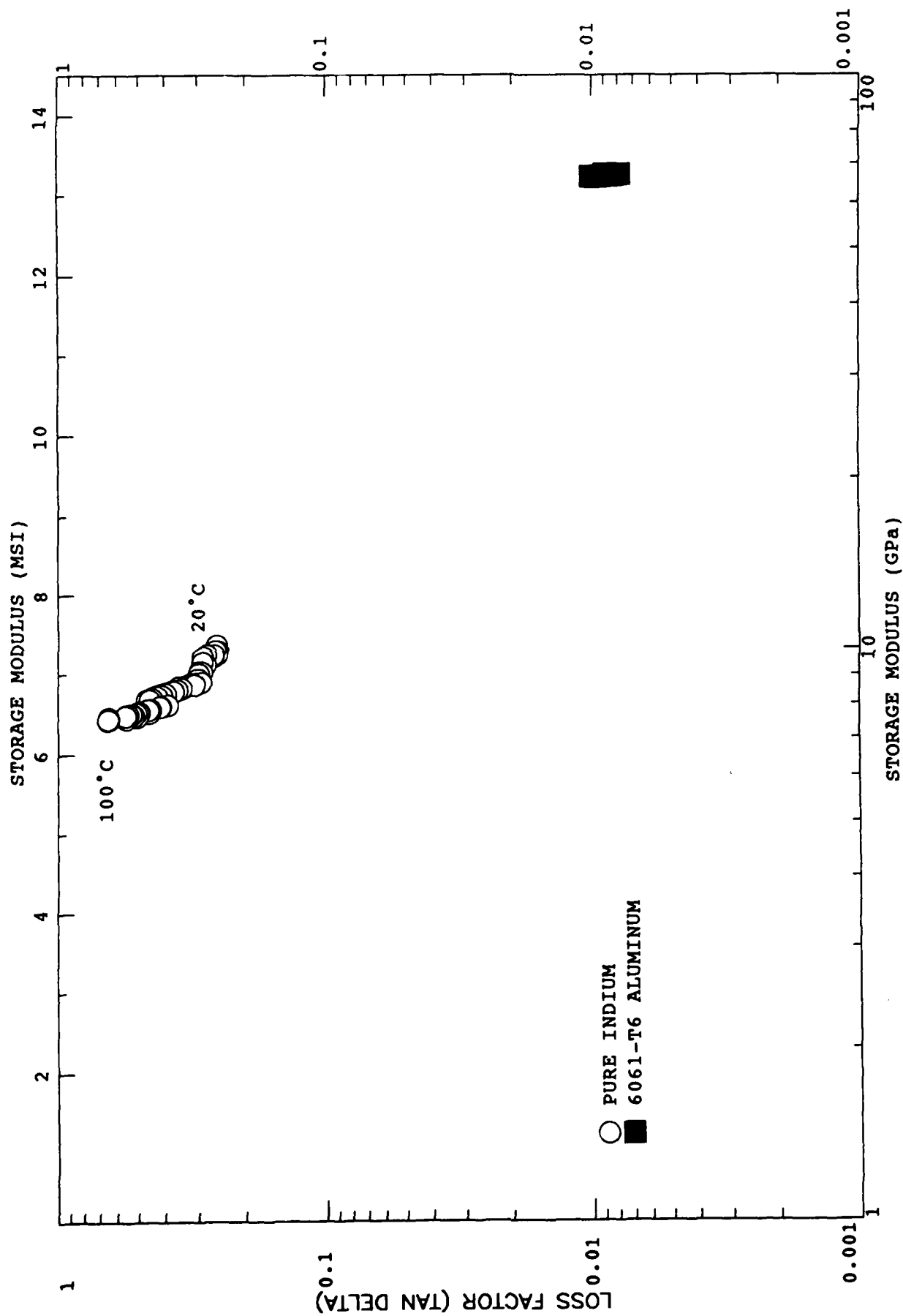


Fig. 4. Values of loss factor and storage modulus of the monolithic materials measured from 20 to 100°C at 0.1 Hz.

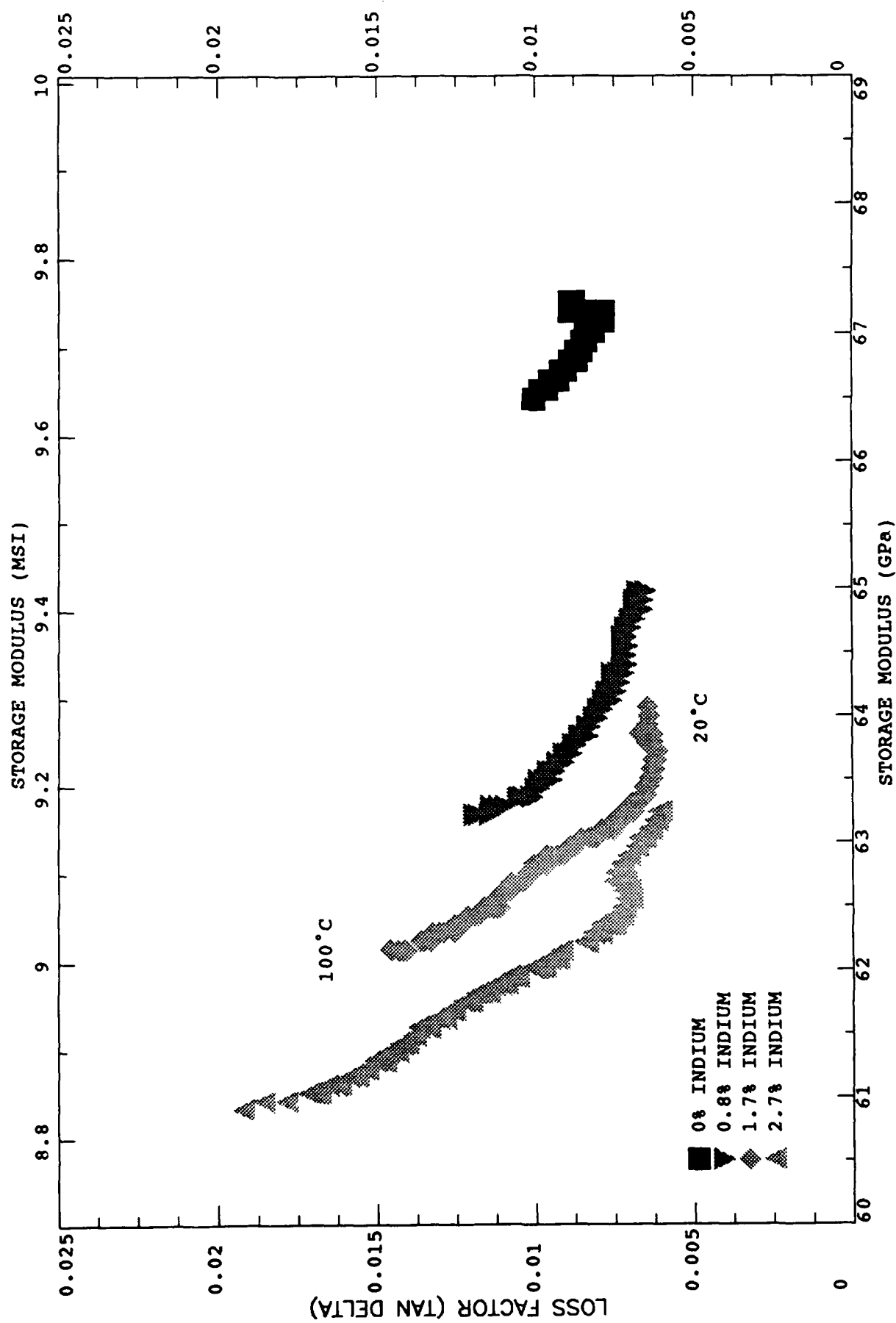


Fig. 5. Values of the loss factor and storage modulus of 6061 aluminum with 0 to 2.7 volume percent indium measured from 20 to 100°C at 0.1 Hz.

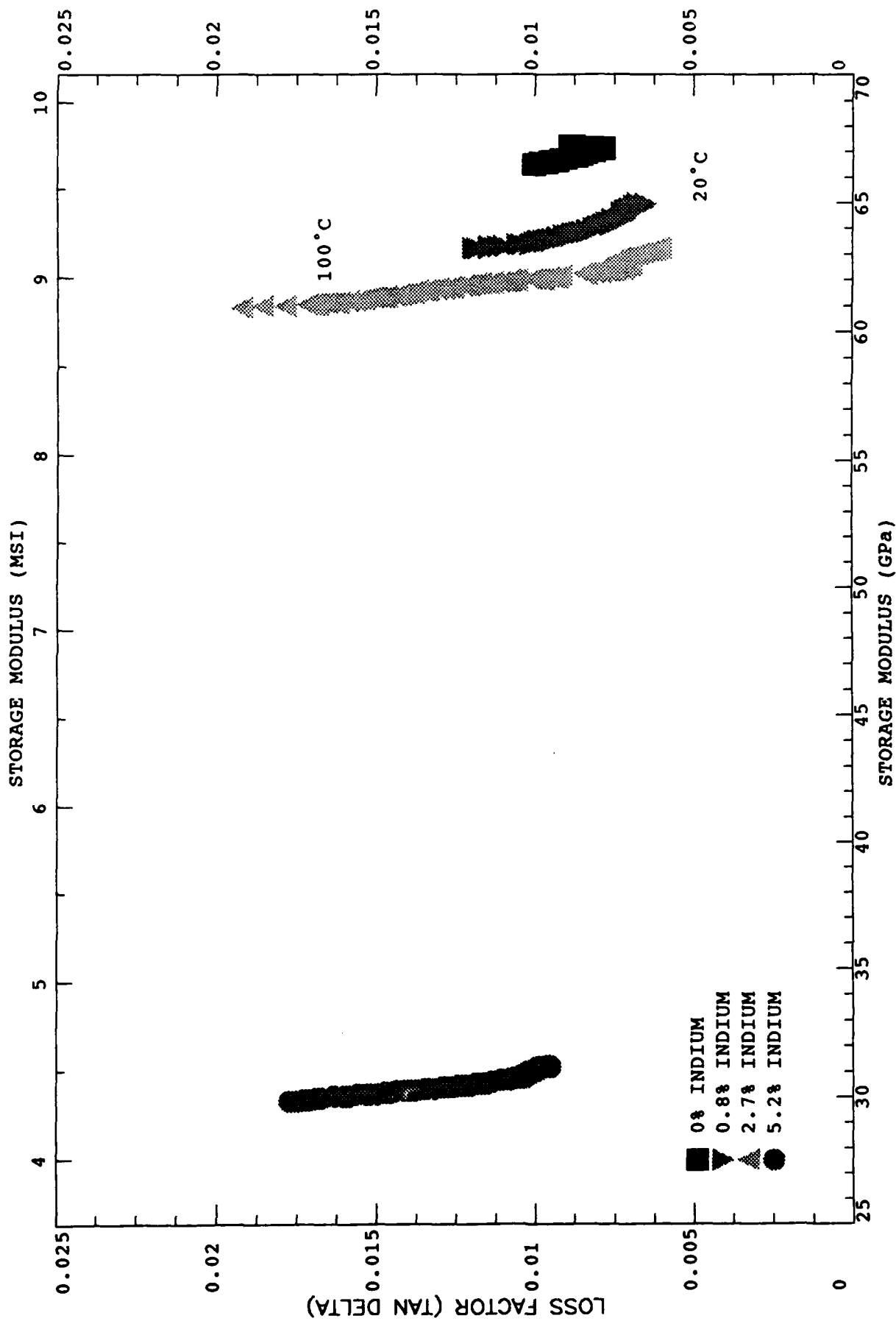


Fig. 6. Values of the loss factor and storage modulus of 6061 aluminum with 0 to 5.2 volume percent indium measured from 20 to 100°C at 0.1 Hz.

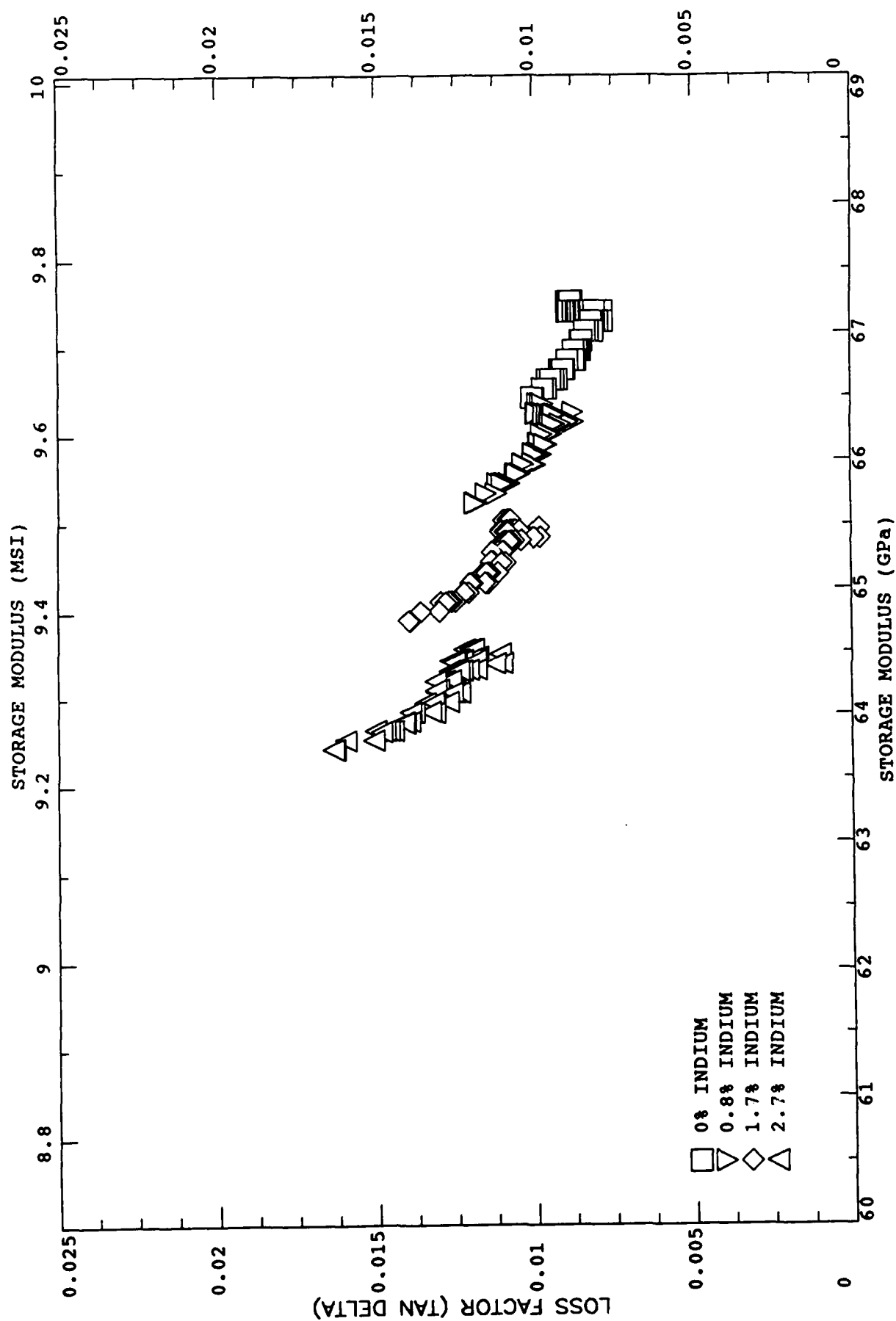


Fig. 7. Values of the loss factor and storage modulus of 6061 aluminum with 0 to 2.7 volume percent indium calculated using Nielsens model.

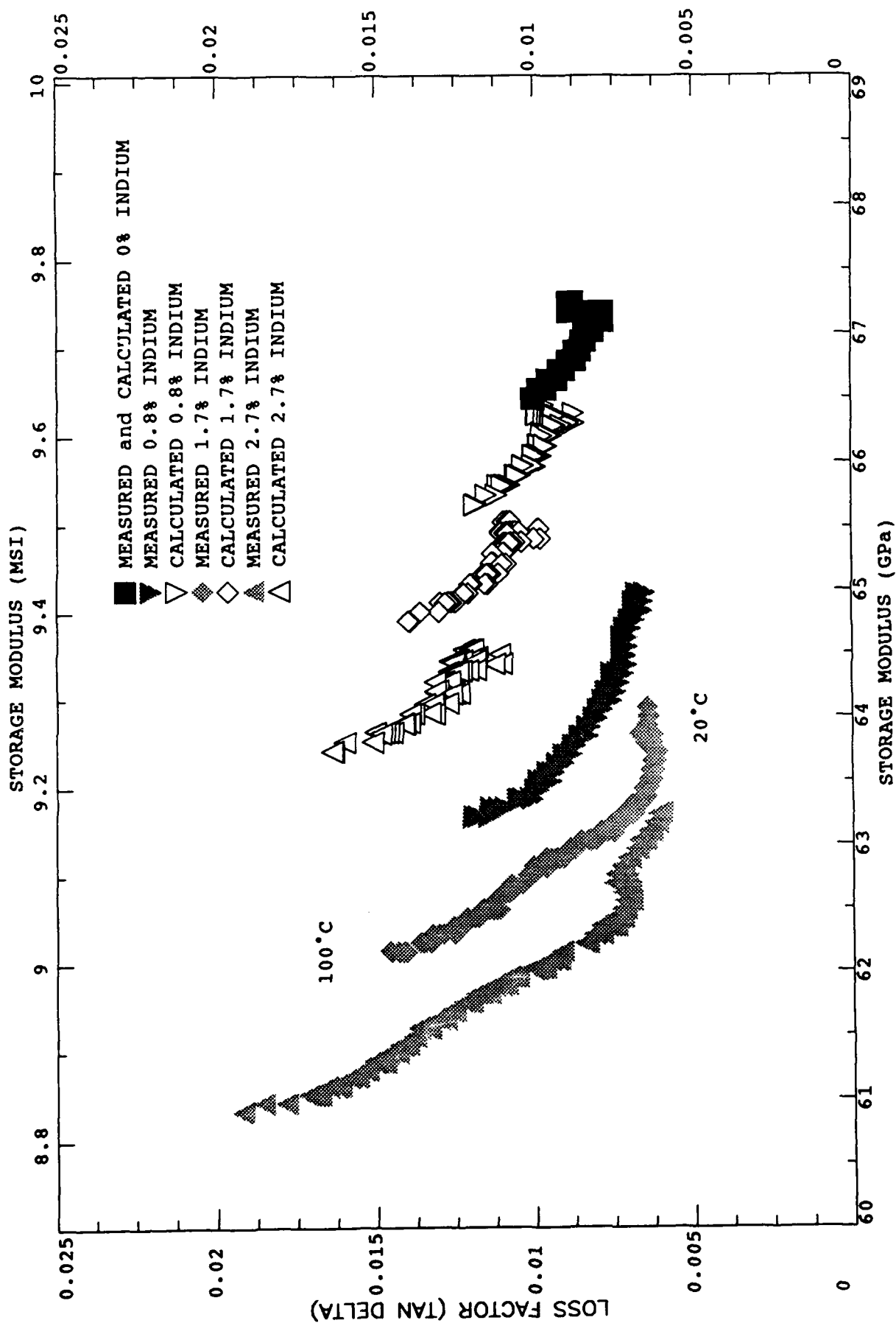


Fig. 8. Comparison of measured and calculated values of the loss factor and storage modulus of 6061 aluminum with 0 to 2.7 volume percent indium.

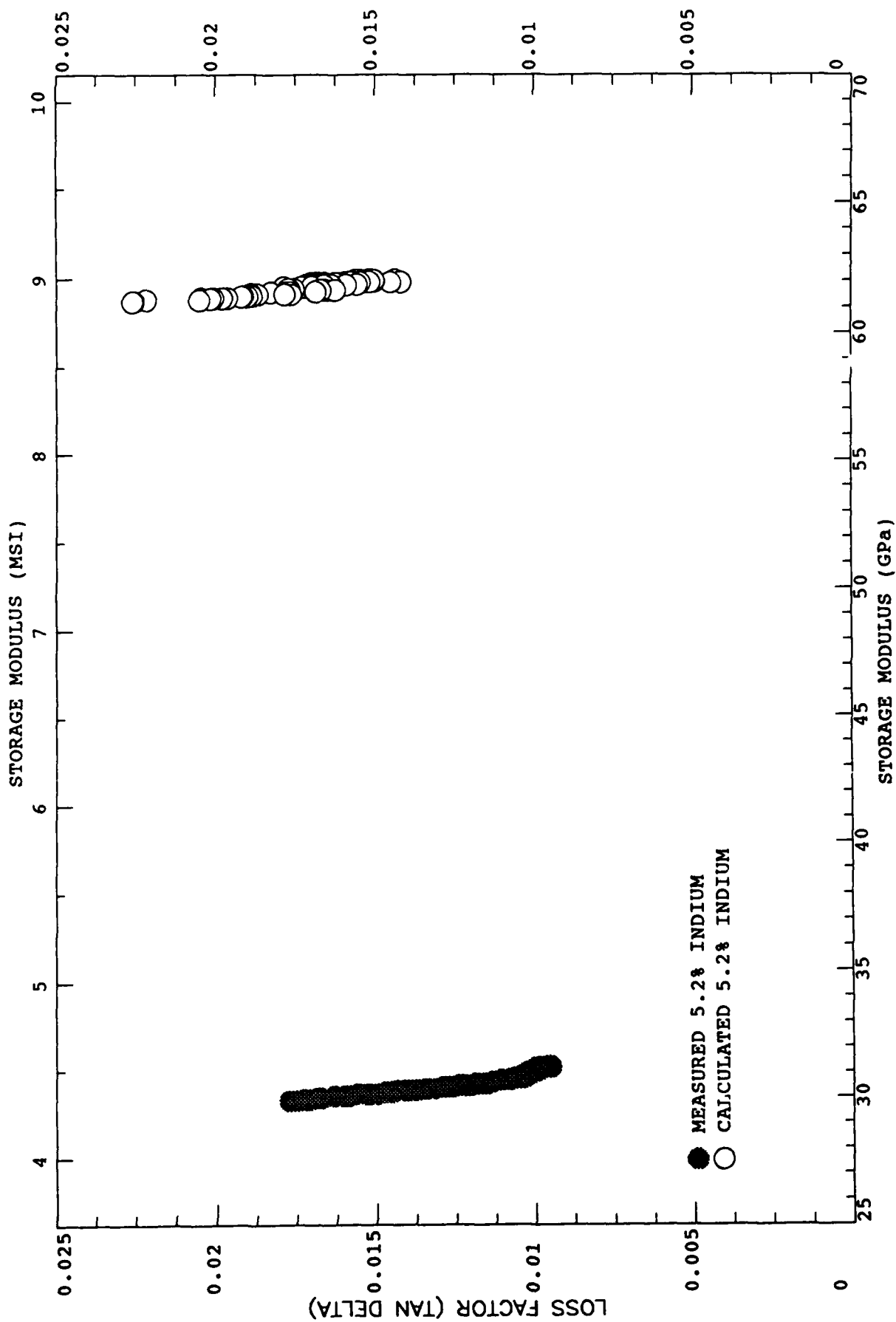


Fig. 9. Comparison of measured and calculated values of the loss factor and storage modulus of 6061 aluminum with 5.2 volume percent indium.

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